# New Electromagnetism V3 The World Leader in Electromagnetic Physics Home of New Electromagnetism WWW. Distinti. Com $K_E = \frac{1}{4}$ 46 Rutland Ave, Fairfield, CT 06825 WWW. Distinti. Com **Ouick Reference Chart**



## Version 3.02 --- 22 May 2007 – Copyright 2007 Robert J. Distinti

**Charge-Force Forms (most fundamental, see ne.pdf)** 

$$\mathbf{F} = \frac{K_{E}Q_{S}Q_{T}\hat{\mathbf{r}}}{\left|\mathbf{r}\right|^{2}}$$
Coulomb's Model (ne.pdf)
$$\mathbf{F} = \frac{K_{M}Q_{S}Q_{T}}{\left|\mathbf{r}\right|^{2}}\left[\left(\mathbf{v}_{T}\bullet\hat{\mathbf{r}}\right)\mathbf{v}_{S} - \left(\mathbf{v}_{S}\bullet\hat{\mathbf{r}}\right)\mathbf{v}_{S} - \left(\mathbf{v}_{S}\bullet\mathbf{v}_{T}\right)\hat{\mathbf{r}}\right]$$

$$\mathbf{F} = \frac{-K_{M}Q_{S}Q_{T}}{\left|\mathbf{r}\right|^{2}}$$
New Induction (ni.pdf)

**Charge-Fragment conversion (ne.pdf)** 

$$\int_{L} \rho_{L} dL = Q \qquad \rho_{L} = dQ / dL$$

$$\int_{L} I d\mathbf{L} = Q\mathbf{v}$$

$$\int_{L} \frac{dI}{dt} d\mathbf{L} = Q\mathbf{a}$$

# **Charge-Field Forms (ne.pdf)**

$$\mathbf{E} = \frac{K_{E}Q_{S}\hat{\mathbf{r}}}{|\mathbf{r}|^{2}} \qquad V_{P} = -\int_{L} \mathbf{E} \cdot d\mathbf{L} \qquad V_{P} = \frac{PE}{Q}$$

$$\mathbf{M} = \frac{K_{M}Q_{S}}{|\mathbf{r}|^{2}} [(\mathbf{v}_{T} \cdot \hat{\mathbf{r}})\mathbf{v}_{S} - (\mathbf{v}_{S} \cdot \hat{\mathbf{r}})\mathbf{v}_{S} - (\mathbf{v}_{S} \cdot \mathbf{v}_{T})\hat{\mathbf{r}}] \qquad V_{K} = \int_{L} \mathbf{M} \cdot d\mathbf{L} \quad \text{For M} \quad V_{K} = \frac{KE}{Q}$$

$$\mathbf{M} = \frac{-K_{M}Q_{S}\mathbf{a}_{S}}{|\mathbf{r}|} \qquad V_{K} = \frac{1}{A} \int \int (\mathbf{M} \cdot d\mathbf{L}) dA \quad \text{For non-uniform M}$$

**Voltage Definitions (ne.pdf)** 

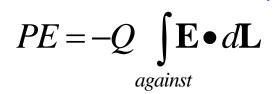
$$V_{P} = -\int_{L} \mathbf{E} \bullet d\mathbf{L} \qquad V_{P} = \frac{PE}{Q}$$

$$V_{K} = \int_{L} \mathbf{M} \bullet d\mathbf{L} \text{ For M uniform in cross-section} \qquad V_{K} = \frac{KE}{Q}$$

$$V_{K} = \frac{1}{A} \iint_{L} (\mathbf{M} \bullet d\mathbf{L}) dA \text{ For non-uniform M}$$

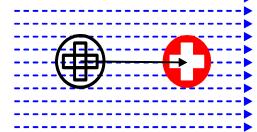
# **Energy Definitions (ne.pdf) (these are still somewhat tentative)**

Q moved against field



$$DE = -Q \int \mathbf{M} \cdot d\mathbf{L} | KE = Q \int \mathbf{M} \cdot d\mathbf{L}$$
against

Q moved by the field



$$KE = Q \int_{with} \mathbf{E} \cdot d\mathbf{L}$$

$$KE = Q \int_{with} \mathbf{M} \cdot d\mathbf{L}$$

The "rule" of nearly instantaneous conversion. This is for low frequency, low loss conditions only (ne.pdf).

$$V_K \cong -V_P$$

# Fragmentary Forms (ne.pdf or nm.pdf)

$$\int_{L} \rho_{L} dL = Q \qquad \rho_{L} = dQ / dL \qquad d^{2}V_{P} = -K_{E} \rho_{LS} \frac{dL_{S} \hat{\mathbf{r}} \cdot d\mathbf{L}_{T}}{|\mathbf{r}|} \quad \text{(ne.pdf)}$$

$$\int_{L} I d\mathbf{L} = Q \mathbf{v}$$

$$\int_{L} \frac{dI}{dt} d\mathbf{L} = Q \mathbf{a}$$

$$d^{2}V_{K} = \frac{-K_{M} I_{S}}{r^{2}} \left[ \frac{I_{S}}{Q_{S}} (d\mathbf{L}_{S} \cdot \hat{\mathbf{r}}) + (\mathbf{v}_{FS} \cdot \hat{\mathbf{r}}) - (\mathbf{v}_{FT} \cdot \hat{\mathbf{r}}) (d\mathbf{L}_{S} \cdot d\mathbf{L}_{T}) + (d\mathbf{L}_{T} \cdot \hat{\mathbf{r}}) (\mathbf{v}_{FT} \cdot d\mathbf{L}_{S}) \right]$$

$$d^{2}V_{K} = \frac{-K_{M} I_{S}}{r^{2}} \left[ \frac{I_{S}}{Q_{S}} (d\mathbf{L}_{S} \cdot \hat{\mathbf{r}}) + (\mathbf{v}_{FS} \cdot \hat{\mathbf{r}}) - (\mathbf{v}_{FT} \cdot \hat{\mathbf{r}}) (\mathbf{v}_{FT} \cdot d\mathbf{L}_{S}) \right]$$

$$d^{2}V_{K} = -K_{M} \frac{dI_{S}}{dt} \frac{d\mathbf{L}_{S} \cdot d\mathbf{L}_{T}}{|\mathbf{r}|} \quad \text{(ne.pdf)}$$

$$d^{2}V_{K} = -K_{M} \frac{dI_{S}}{dt} \frac{d\mathbf{L}_{S} \cdot d\mathbf{L}_{T}}{|\mathbf{r}|} \quad \text{(ne.pdf)}$$

(\*) = unfinished or unreleased (!)=tentative, requires more investigation

Wire Forms (ne.pdf). These are derived by integrating the fragmentary forms. For simplicity, these assume uniformity over the cross-section which may not be applicable in some cases. These integrations can be further simplified by assuming that  $\rho$  and I are uniform along the length of the wire.

$$V_{P} = -K_{E} \int_{S} \rho_{LS} \frac{dL_{S} \hat{\mathbf{r}} \cdot d\mathbf{L}_{T}}{|\mathbf{r}|^{2}} \quad \text{(ne.pdf)}$$

$$V_{K} = -K_{M} \oint_{S} \frac{I_{S}}{|\mathbf{r}|^{2}} \left[ \frac{((\mathbf{v}_{FS} - \mathbf{v}_{FT}) \cdot \hat{\mathbf{r}})(d\mathbf{L}_{S} \cdot d\mathbf{L}_{T})}{+(d\mathbf{L}_{S} \cdot \hat{\mathbf{r}})(\mathbf{v}_{FS} \cdot d\mathbf{L}_{T})} + (d\mathbf{L}_{T} \cdot \hat{\mathbf{r}})(\mathbf{v}_{FT} \cdot d\mathbf{L}_{S}) \right] \quad \text{(nm.pdf)}$$

$$V_K \cong -V_P \qquad V_K = -K_M \int_S \int_T \frac{dI_S}{dt} \frac{d\mathbf{L}_S \bullet d\mathbf{L}_T}{|\mathbf{r}|} \quad \text{(ne.pdf)}$$

New Electromagnetism represents a more complete accounting of electromagnetic interactions than classical theory (Maxwell et al). These models should never be considered "finished." As more phenomena are encountered or as conventions are updated, these models will be revised accordingly.

Present changes under consideration include the use of "I-Field" for New Induction instead of M-Field. This change may affect the energy definitions.

All pdf files referenced can be found at www.distinti.com/docs

**Kirchoff's rules and curl:** These are valid for both static and dynamic conditions.

$$\oint_{L} \mathbf{E} \cdot d\mathbf{L} = 0$$

$$\nabla \times \mathbf{E} = 0$$

$$\oint_{L} \mathbf{M} \cdot d\mathbf{L} = V_{K}$$

$$\nabla \times \mathbf{M} = (*)$$

#### Poisson

$$\nabla^2 V_P = -4\pi K_E \rho_V$$
$$\nabla^2 V_K = 0(!)$$

### **Divergence**

$$\nabla \bullet \mathbf{E} = 4\pi K_E \rho_V$$
$$\nabla \bullet \mathbf{M} = 0(!)$$

### Gradient

$$egin{array}{c} 
abla V_P = -\mathbf{E} \\ 
abla V_K = (*) 
onumber \end{array}$$

The B and H fields of classical theory, as well as all derived constructs, are obsolete. See the secrets of qvxb.pdf for details.